

High-Sensitivity Instantaneous Microwave Frequency Measurement Based on a Silicon Photonic Integrated Fano Resonator

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(Invited Paper)

Abstract—Instantaneous frequency measurement (IFM) based on a silicon photonic Fano resonator with improved linearity and sensitivity is proposed and experimentally demonstrated. The Fano resonator has a steep edge in its spectral response, which is employed to translate the frequency information of a microwave signal to an optical power change. When comparing the optical powers at the output and input of the Fano resonator, a highly linear amplitude comparison function (ACF), which is used to estimate the microwave frequency is obtained. The key device in the system is the Fano resonator, which is realized by coupling a grating-based Fabry–Perot cavity resonant mode with an add-drop microring resonator mode, implemented on a silicon platform. The linearity of the ACF is characterized by its R-squared value which is calculated by fitting the ACF measurements with a linear function. In our experimental demonstration, an R-squared value as large as 0.99 is obtained. A frequency measurement range as large as 15 GHz with a resolution better than ± 0.5 GHz is achieved. The use of the proposed IFM system to perform Brillouin frequency discrimination in a fiber-optic sensor for temperature measurement is demonstrated.

Index Terms—Carrier-suppressed single sideband modulation, Fano resonance, instantaneous frequency measurements, microwave photonics, silicon photonics.

I. INTRODUCTION

INSTANTANEOUS frequency measurement (IFM) of a microwave or millimeter-wave signal is of great importance

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for applications in electronic warfare and cognitive radio systems and for high resolution sensor systems. In recent years, photonics-based techniques for IFM have been intensively investigated and numerous approaches have been proposed and demonstrated. The key advantages of performing IFM based on photonics include ultra-fast measurement speed, high accuracy, and large frequency measurement range [1], [2]. In general, photonics-based IFM techniques can be divided into three categories: frequency-to-microwave power mapping [3]–[8], frequency-to-optical power mapping [9]–[22], and frequency-to-time mapping [23]–[25].

For techniques based on frequency-to-microwave power mapping [3]–[8], a microwave signal with its frequency to be measured is first converted into an optical signal through electrical to optical conversion using an optical modulator, and then sent the modulated optical signal to a photonic processing module (a dispersive element or an optical filter), where a monotonic frequency-dependent microwave power penalty function in the optical domain is introduced. A photodetector (PD) is connected to the photonic processing module to convert the optical signal into a microwave signal. To make the measurement to be independent of the input power, two parallel channels which provide different, usually opposite frequency-dependent power penalty functions, are needed. The power ratio of the microwave signals from the two channels, called amplitude comparison function (ACF), is a function of the input microwave frequency. The major problem associated with these techniques is that high-frequency PDs and wideband microwave devices are needed, making the systems costly.

To avoid the use of expensive microwave devices, techniques based on frequency-to-optical power mapping were proposed and the ACF is the power ratio between the optical powers measured at the outputs of two optical paths [9]–[22]. The key component in such a system is a photonic processing module, such as an optical filter [9]–[16] or an optical mixing unit [17]–[22], to translate the frequency of an input microwave signal to an optical power change. In most of the IFM schemes based on frequency-to-optical power mapping, an optical filter with a sinusoidal spectral response is employed to perform frequency to power mapping, which leads to an ACF with very

poor linearity. For example, a Sagnac loop based filter with 50 GHz free spectral range (FSR) was used in [9]. The microwave signal with an unknown frequency was modulated on two optical carriers, one of which was set at a peak and the other at a valley of the filter spectral response. The power ratio of the two modulated signals, which is the ACF, is given by [9]

$$\text{ACF} = \frac{1 - \cos(2\pi f_{RF}/\text{FSR})}{1 + \cos(2\pi f_{RF}/\text{FSR})} \quad (1)$$

where f_{RF} is the frequency to be measured. The R-square value is calculated to be 0.65806, which is very low, indicating a poor linearity.

For techniques based on frequency-to-time mapping [23]–[25], the frequency of an unknown microwave signal is converted into an electrical time delay using a dispersive time delay device [23] or a frequency shifting recirculating delay line (FS-RDL) [24], [25]. Since a dispersive element has a linear group delay and a FS-RDL has equal frequency shift per roundtrip time, these techniques can lead to a linear IFM system. However, the dispersion-based method has a relatively low resolution (~ 12.5 GHz) and a large measurement error (about ± 1.56 GHz). The FS-RDL-based method has a long response time, which may not be suitable for instantaneous frequency measurements.

To achieve an IFM system with a high linearity, high resolution, and short response time, a photonic processing module with a highly linear frequency response to perform frequency to power conversion is required. A linear ACF would ensure identical measurement sensitivity for an entire frequency measurement range, which is important for an IFM system to have a high and uniform measurement resolution and accuracy. For instance, a special FBG with two linear and complementary spectral slopes was applied to achieve a highly linear ACF [13]. A measurement range of 1–10 GHz and a measurement accuracy of ± 0.2 GHz were demonstrated [13].

Among the approaches [3]–[25], those reported in [7], [8], [14], [15], [18], [19], [22] were implemented based on integrated photonic devices. The size is smaller, the stability is better, the measurement latency is shorter, and the unit cost is lower. These make integrated IFM attractive for practical applications. However, those integrated solutions have poor linearity, making the measurement resolution not uniform over the frequency measurement range.

The performance of the integrated IFM techniques is summarized in Table I, in which the measurement range, measurement error and input power to the chip are considered. As can be seen the input power levels for all the approaches [7], [8], [14], [15], [18], [19], [22] were very high. For a practical IFM system, the input power from a receiver is usually very low. Thus, a solution to perform IFM with a lower input power is needed.

In this paper, we propose and experimentally demonstrate a novel approach to achieve IFM based on a silicon photonic integrated Fano resonator with high linearity, high sensitivity and low input power. The key device in the system is the Fano resonator which is implemented using a grating-based

TABLE I
PERFORMANCE OF IFM TECHNIQUES BASED ON INTEGRATED PHOTONICS

Techniques	Measurement range (GHz)	Measurement error (GHz)	Input Power (dBm)
SBS [7]	0-38	± 0.001	20
FBG [8]	0-32	± 2.5	10
PM-IM conversion [14]	0.5-4	± 0.094	20
Ring-assisted MZI [15]	5-15	± 0.2	12
Nonlinear effect [18]	$> 2.5 \times 10^3$	-	22
Nonlinear effect [19]	640 Gbit/s	-	14
Nonlinear effect [22]	0-40	± 0.318	> 20

Fabry-Perot (FP) cavity-coupled add-drop microring resonator (MRR), where the resonant mode of the FP cavity interferes with the resonant mode of the MRR in transmission response to produce the Fano resonance. When an optical wave that carries a microwave signal is applied to the Fano resonator, the frequency information is converted to an optical power. By calculating the ratio between the optical powers at the input and the output ports of the Fano resonator, an ACF is obtained and the frequency of the microwave signal is measured. A proof-of-concept experiment is carried out. With only 0 dBm input power to the chip, a frequency measurement range as large as 15 GHz with a resolution better than ± 0.5 GHz is achieved. The use of the proposed approach for Brillouin frequency discrimination in a fiber-optic sensor for high resolution temperature measurement is studied and demonstrated. A measurement sensitivity as large as 1.24 MHz/ $^{\circ}$ C is achieved.

II. PRINCIPLE

The schematic of the proposed system using a Fano resonator for instantaneous microwave frequency measurement is shown in Fig. 1. A continuous wave (CW) optical carrier (λ_c) is modulated by a microwave signal with its frequency to be measured at a Mach-Zehnder modulator (MZM) that is biased at the minimum transmission point (MITP) to generate a double-sideband modulated signal with suppressed carrier. After filtering out the lower sideband using an optical bandpass filter (OBPF), a carrier-suppressed single-sideband (CS-SSB) signal is achieved and launched into the silicon chip via an optical circulator (OC). Part of the CS-SSB signal is transmitted through the Fano resonator while the other part of the CS-SSB signal is reflected by the grating. The ACF, defined as the power ratio between the transmitted and reflected average optical powers, is given

$$\text{ACF}(\lambda) = \frac{P_{\text{TX}}(\lambda)}{P_{\text{RX}}(\lambda)} \quad (2)$$

Since the frequency of the microwave signal has a unique relationship with the ACF, once the ACF value is known, the frequency of the microwave signal is measured. Note that the optical powers rather than the microwave powers are measured, thus no high speed PDs are needed.

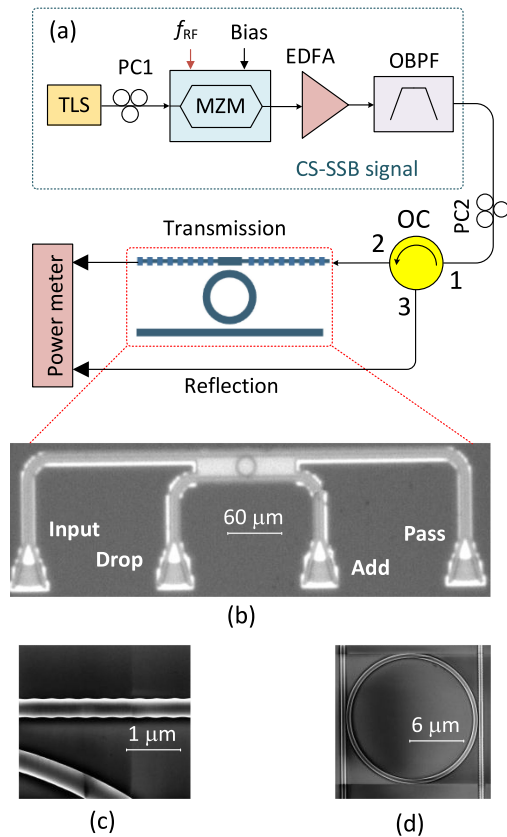


Fig. 1. (a) Schematic diagram of the proposed IFM system. (b) Image of the fabricated Fano resonator obtained from a microscope camera. (c) The SEM image of the Bragg grating. (d) The SEM image of the MRR. CS-SSB: carrier-suppressed single-sideband; TLS: tunable laser source; PC: polarization controller; MZM: Mach-Zehnder modulator; EDFA: erbium-doped fiber amplifier; OBPF: optical bandpass filter; OC: optical circulator.

III. FREQUENCY MEASUREMENT OF A MICROWAVE SIGNAL

A proof-of-concept experiment is carried out based on the setup in Fig. 1. An optical carrier emitting at 1556.2 nm with a power of 8 dBm is generated from a tunable laser source (TLS, Anritsu MG9638A), which is then modulated by a microwave signal from a microwave source (Agilent E8254A) at a 20-GHz MZM (JDS Uniphase, Model 10026465). The MZM is biased at the MITP and the microwave signal is amplified by an RF amplifier (MultiLink modulator driver MTC5515-751) by 20 dB. Then, the carrier-suppressed double-sideband (CS-DSB) signal is amplified using an erbium-doped fiber amplifier (EDFA, Norstel Inc.) before introduced into the OBPF (Finisar, Waveshaper 4000 s). The lower sideband (LSB) is selected and launched into the chip and the input power is kept around 0 dBm since the EDFA works in a constant output power mode (Automatic Power Control, APC). The total loss of the chip is 13.6 dB. The optical powers at the transmission and reflection ports are simultaneously measured by a two-channel optical power meter (HP 8152A). Since the Fano resonator will experience a spectral shift when the temperature changes, accurate and stable temperature control (TEC, ILX Lightwave LDT-5910B) is employed to make the temperature of the chip fixed at 23 °C. Thanks to

the stable environmental temperature and low input power to the chip, the TEC controller can easily make the temperature stabilized with a temperature variation smaller than 0.01 °C during the process of the experiment and the power ratio of the two channels is calculated and recorded. Since the two parallel channels are experiencing the same environmental fluctuations, the impact of the temperature change to the ACF is small and can be neglected.

The transmission and reflection spectral responses of the Fano resonator, which was first reported by us in [26] are measured by an optical vector analyzer (LUNA OVA CTe) and shown in Fig. 2(a) and (b). A zoom-in view of the region around the Fano resonance is shown in Fig. 2(c) and (d). The ACF is calculated using (2), which is shown in Fig. 2(e). A zoom-in view of the linear region of the ACF is shown in Fig. 2(f) covering a frequency range of 15 GHz. The R-square for the linear region is calculated which is 0.98658.

Fig. 3 shows the measured microwave frequencies when the output powers from the microwave source are set at -5 dBm, 0 dBm and 5 dBm. The frequency tuning range is from 3 to 18 GHz, which cannot be higher due to the limited bandwidth of the 20 GHz MZM. The solid line shows the actual ACF of the system. The measurement errors which are the difference between the measured ACF values and the actual ACF values are shown in Fig. 4. The measurement errors are smaller than ± 0.5 GHz over the entire frequency measurement range from 3 to 18 GHz. The standard deviations of the errors for the three power levels are calculated to be 0.1331, 0.1227, and 0.1625 (0.1142 without counting the first two points which have higher errors). As can be seen, a higher input power will lead to smaller measurement errors. A higher microwave power results in a higher modulation index so that the signal-to-noise ratio (SNR) of the modulated optical signal is higher after amplification. The measurement errors for lower frequencies are higher when the signal power is 5 dBm. This is because the high order sidebands are stronger which are harder to be removed by the optical bandpass filter.

The approach is effective for the measurement of a single-frequency microwave signal. If a modulated microwave signal is applied to the input of the IFM system, each frequency component of the input microwave signal will be mapped to a corresponding optical power and the output power will be the sum of all the optical powers. Since an R-squared value as large as 0.99 is obtained in the proposed IFM system, the power-frequency relation is highly linear and the measurement result will be a linear combination of the frequencies, which will cause measurement errors. Taking a double-sideband plus carrier (DSB+C) signal as an example, if the carrier frequency is 10 GHz, and each of the two sidebands has a power that is 5% or 10% of the power of the microwave carrier, we calculate the measurement errors as a function of the sideband frequency, as shown in Fig. 5. As can be seen, the measurement error decreases as the frequency of the sidebands and the sideband-to-carrier power ratio decrease. If the two sidebands are separated from the microwave carrier by 1 GHz, the frequency measurement error is 0.2 GHz (or 2%) or 0.4 GHz (or 4%).

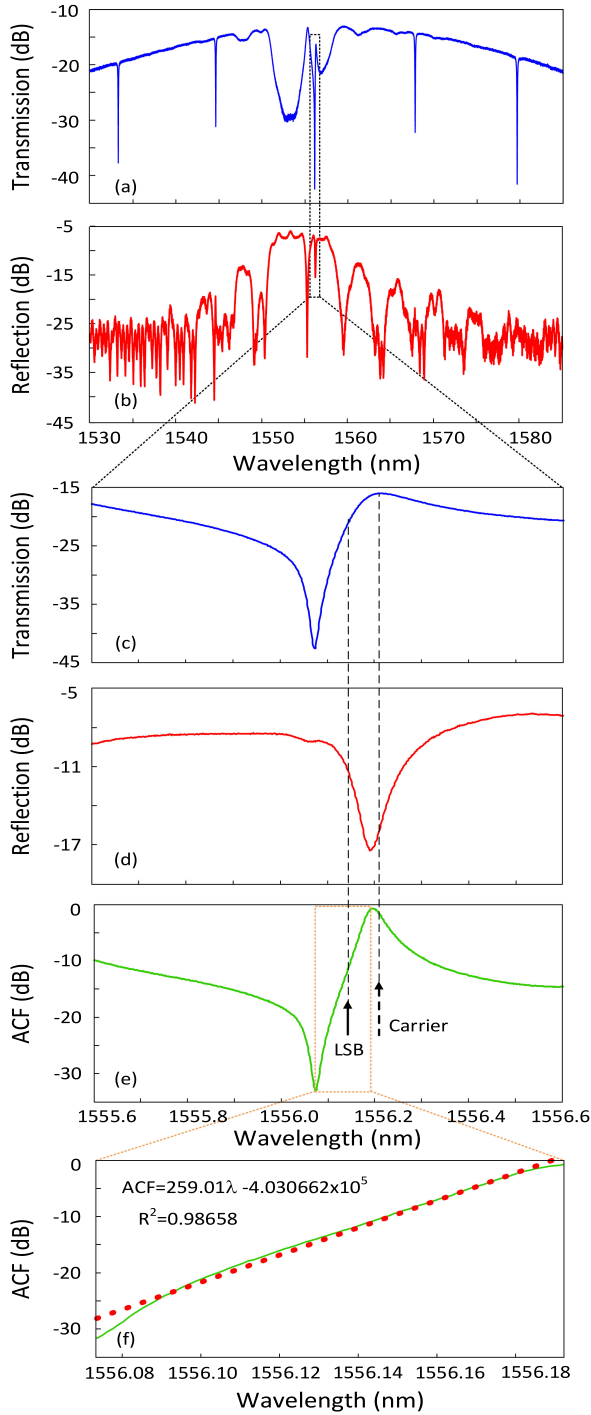


Fig. 2. (a) Transmission spectral response of the device, (b) reflection spectral response of the device, (c) and (d) zoom-in views of the region around the Fano resonance, (e) ACF, and (f) a zoom-in view of the linear region of the ACF with linear fitting.

IV. FREQUENCY DISCRIMINATION IN A BRILLOUIN FIBER SENSING SYSTEM

Distributed Brillouin-based optical fiber sensors are widely used in civil infrastructure and power supply systems to monitor the temperature and strain variations along the sensing fibers [27]. High speed and high resolution interrogation of

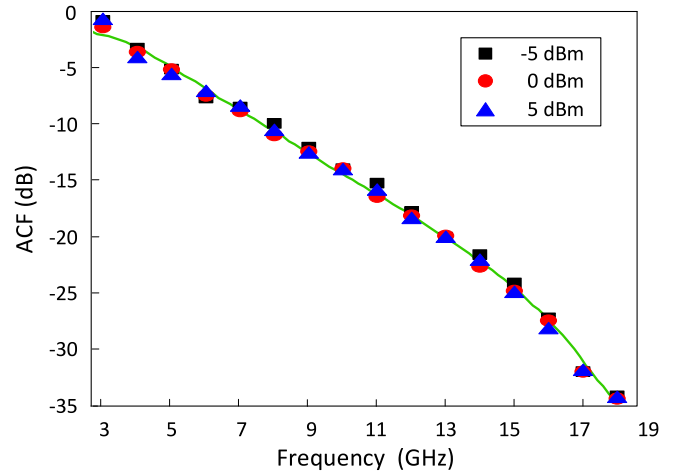


Fig. 3. Measured microwave frequencies when the input powers of the unknown microwave signal are -5 dBm, 0 dBm and 5 dBm.

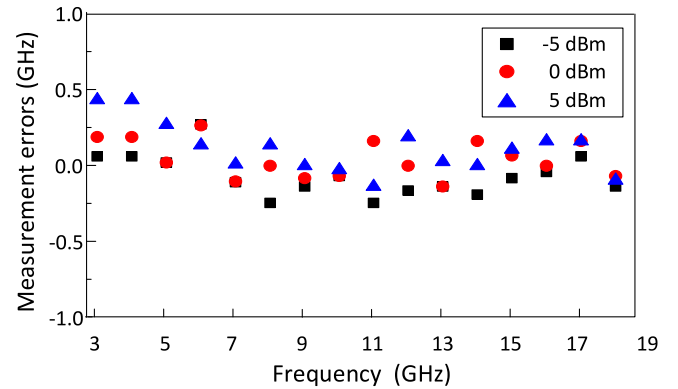


Fig. 4. Measurement errors when the input powers of the microwave signal are at -5 dBm, 0 dBm and 5 dBm.

the Brillouin frequency shift (BFS) of such a sensor is a key procedure to ensure a high performance temperature and strain sensing [28], [29]. Thanks to the high speed and high linearity, the proposed IFM system is a good candidate to achieve real-time BFS discrimination in a Brillouin fiber-optic sensing system.

Fig. 6 shows the experimental setup to use the proposed IFM technique to interrogate a Brillouin sensor. A CW light emitted from a TLS (Anritsu MG9638A) is amplified by an EDFA (Nortel) to 15.8 dBm and then launched into a 10 -km single-mode fiber (SMF) through an optical circulator (OC1) as a pump source. Due to stimulated Brillouin scattering, a Brillouin gain spectrum at a longer wavelength is resulted, as shown in Fig. 7. A peak at the same wavelength as the pump wavelength due to Rayleigh scattering is seen, which is removed by an OBPf (Finisar, Waveshaper 4000s), as also shown in Fig. 7. The Brillouin gain spectrum is introduced into the IFM chip through a second OC (OC2). By detecting the transmission and reflection powers, the optical frequency of the Brillouin gain spectrum is estimated. To evaluate its performance as a temperature sensor, in the experiment we change the temperature of the SMF from 75 °C to 25 °C, the values of the BFS

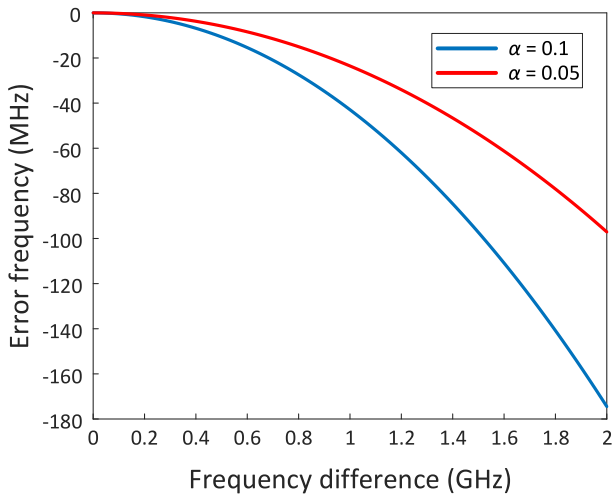


Fig. 5. Measurement errors of a DSB+C signal with 10 GHz microwave carrier when the sideband-to-carrier power ratios are 5% and 10%.

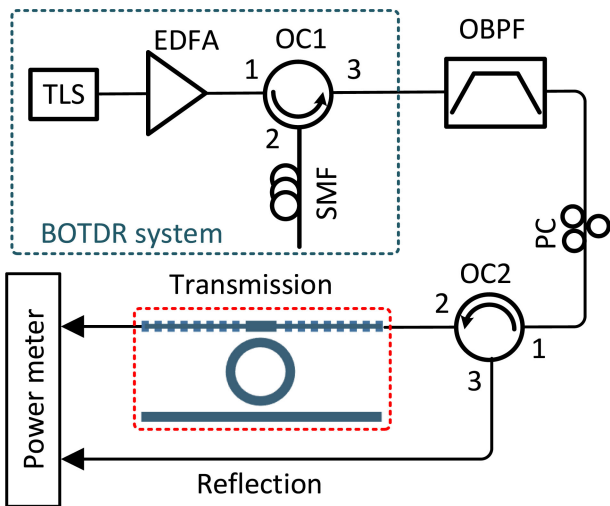


Fig. 6. Experiment setup of the proposed microwave frequency measurement system. BOTDR: Brillouin optical time domain reflectometry; SMF: single-mode fiber.

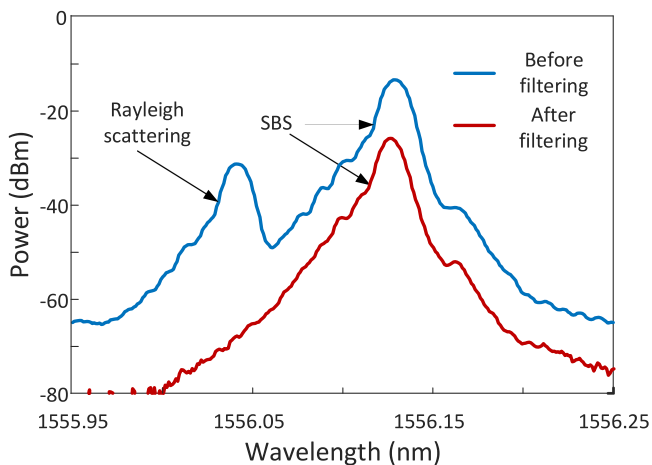


Fig. 7. Optical spectra before and after optical filtering.

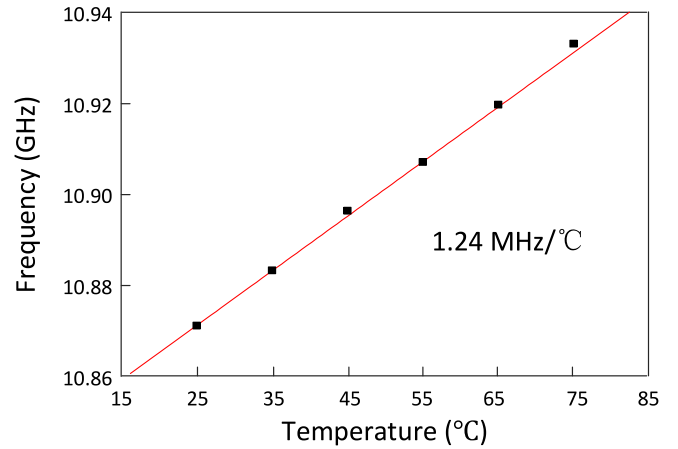


Fig. 8. Brillouin frequency shift as a function of the temperature change of the SMF.

are measured. As can be seen from Fig. 8 the values are decreased with a fitted slope of 1.24 MHz/°C. The measurement errors are within ± 1.64 MHz, which are mainly caused by the nonuniformity of the temperature distribution in such a long fiber as well as the strain variations caused by the temperature change.

Note that if the CW light source and the two-channel optical power meter are replaced by an optical pulse source and two low-speed PDs, the sensor can be used for high speed and high resolution distributed temperature sensing. For practical implementations, temperature control of the chip can be used to control the temperature of the chip in real time so that the measurement errors would be further reduced.

V. CONCLUSION

A novel approach to implement instantaneous frequency measurement based on a silicon photonic Fano resonator with high linearity, high sensitivity and low input power was proposed and experimentally demonstrated, for the first time, to the best of our knowledge. The Fano resonance was generated by coupling a grating-based FP cavity with an MRR. When the resonant mode of the FP cavity interfered with the resonant mode of the MRR, an asymmetric Fano response with a steep edge in its spectral response was realized, which was used to convert the frequency of a microwave signal to an optical power change. By comparing the optical powers at the output and the input ports of the Fano resonator, an ACF with high linearity was obtained. An experiment was performed. The R-squared value of the measured ACF was 0.99 while is very high confirming a good linearity of the ACF. The input power to the chip was about 0 dBm, which was low. A frequency measurement range as large as 15 GHz with a resolution better than ± 0.5 GHz was achieved. The proposed IFM system was also used for Brillouin frequency discrimination in a fiber-optic sensor for temperature measurement. A linear dependence between the Brillouin frequency shift and the temperature variations was measured to be 1.24 MHz/°C with measurement errors within ± 1.64 MHz.

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